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sales@silixa.com
Missoula, MT: +1 406 204 7298
Houston, TX: +1 832 772 3333
UK: +44 208 327 4210
Silixa.com
Twitter
LinkedIn
YouTube
Distributed sensing enables continuous, real-time measurements along a length of optical fibre.

Advances in optoelectronics and associated signal processing have enabled the development of optical fibre distributed sensors with maximum ranges of several tens of kilometres. Unlike traditional sensing that relies on discrete sensors measuring at pre-determined points, distributed sensing utilises the optical fibre as the sensing element without any additional transducers in the optical path.

An optical fibre distributed sensor emits pulses of laser light into an optical fibre. The emitted laser light is scattered within the fibre which causes both attenuation of the pulse amplitude with distance travelled in the fibre and backscattered light that is detected and used to derive measurements of different physical properties along the optical fibre. The position of scattering elements along the fibre is determined from the time of flight between the emission of a laser pulse and the detection of backscattered light through application of the principles of Optical Time-Domain Reflectometry (OTDR). Several scattering processes take place when the pulse of laser light interacts with the molecules of the optical fibre and different measurements can be derived from analyses of the full spectrum of detected light.

Most of the emitted light is backscattered without experiencing a change in wavelength through elastic Rayleigh scattering. True distributed acoustic sensors (DAS) use the Rayleigh scattering signal to derive the coherent full acoustic field (amplitude, wave-
length, and phase) over a wide dynamic range allowing for characterisation of localised acoustic or seismic environments.

Distributed strain and temperature sensors (DSTS) use the interaction of emitted light with lower-frequency molecular vibrations (also referred to as material waves) within a fibre, known as Brillouin scattering, to derive the distribution of coupled strain and temperature across the entire length of a fibre. Distributed temperature sensors (DTS) however make use of wavelength-shifted backscattered light caused by inelastic interactions between the source light and temperature-dependent molecular vibrations within the fibre, known as Raman scattering, to determine strain independent absolute temperature.

DTS, DSTS, and DAS systems have seen widespread use in the oil and natural gas industry to assist reservoir engineers in characterising assets and optimising production and the lifetime of wells. Today, distributed sensing offers integrity and surveillance monitoring solutions for pipelines, water infrastructure such as dams and dikes, dynamic rating of power cables, and process plants. These systems also offer an unprecedented opportunity to obtain high-resolution, real-time temperature measurements in an increasingly wide variety of applications. Distributed fibre optic monitoring offers dense spatial and temporal profiling over large surfaces, long lengths, and at locations where conventional point sensing is not applicable or cost effective.

In this document, the principles of DTS will be introduced. DTS have been proven to be the instruments of choice for numerous environmental applications, including hydrologic applications such as groundwater-surface water interactions in streams and flow dynamics in boreholes. Applications in the atmospheric sciences and oceanography are becoming more commonplace, as are installations to monitor sensitive environments such as the cryosphere.

A DTS system in concert with active heating can also be used to characterise thermal properties and related processes in addition to characterisation based on measuring ambient temperature. Example applications of active DTS (A-DTS) systems include determination of soil moisture, characterisation of apparent thermal conductivity, and detection and quantification of aquifer flow in fractures and estimation of groundwater velocity.
The underlying principle of distributed temperature sensing is a Raman scattering-based temperature measurement combined with optical time-domain reflectometry.

To generate back-scattered light, a DTS unit launches a short pulse of light into an optical fibre. The forward propagating light generates Raman backscattered light at two new wavelengths, from all points along the fibre (FIGURE 1). The wavelengths of the Raman backscattered light differ from the forward propagating light and are named Stokes and anti-Stokes.

The amplitudes of the Stokes and anti-Stokes light are monitored by the DTS unit and the spatial localisation of the backscattered light is determined through knowledge of the propagation speed inside the fibre. The determination of the source of a light signal by measuring the time between the injection of a light source and the detection of a backscattered signal is the fundamental principle of OTDR. The amplitude of the Stokes light is very weakly dependent on temperature, while the amplitude of the anti-Stokes light is strongly dependent on temperature. The temperature at each sampling location is calculated by taking the ratio of the amplitudes of the measured anti-Stokes and Stokes light.

Raman is the weakest of the backscattering effects. Therefore, backscattered light from many pulses must be collected and stacked to get a useful signal from which temperature can be derived above the noise of the system. The laser pulse frequency is in the kHz range for most DTS systems. Note that as the distance to be sensed increases the pulse frequency must decrease to account for the longer time it takes light to travel in the optical fibre; a new pulse cannot be launched until the backscattered light from the previous pulse has been detected (i.e., light from two pulses cannot be present in the fibre at the same time).

FIGURE 1
The DTS unit emits pulses of laser light into an optical fibre and detects the backscattered light. The temperature along the fibre can be determined from the time between the emission and detection of light and the ratio of the amplitudes of backscattered light from specific locations along the fibre.
The key components of DTS systems

**A pulsed laser**

The laser wavelength and pulse repetition rate must be optimised for the target measurement range (FIGURE 2). Wavelengths of light that experience low attenuation in optical fibre (e.g. 1550 nm) are more suitable for long range systems. Wavelengths of light that experience higher attenuation in optical fibre are optimal for shorter ranges (e.g. 1064 nm) because a stronger backscattered light signal is generated. The pulse width, or duration, must be optimised for the target spatial resolution (shorter pulse widths provide better spatial resolution).

**Optoelectronic detectors**

The gain and bandwidth of detectors used to measure the backscattered light signals must be optimised for the target Signal-to-Noise (SNR) and spatial resolution.

**Data acquisition card**

The system must be capable of spatially distributing the backscattered signal by accounting for the time of flight of the light and the sampling frequency must be specified for the target spatial resolution. The linearity of the signal digitisation circuitry is essential for accurate temperature estimation.

**Optical fibre**

Fibre properties must be conducive to meaningful DTS measurements with the chosen system configuration and the target application.
In distributed sensing, optical fibres act as both the waveguide for transmitting pulses of light and as the sensor itself.

Fibre optic sensors, like DTS, that use the fibre directly for measurement are known as intrinsic fibre optic sensors since no other components are needed along the optical fibre. Optical fibre selection is therefore an important consideration in distributed sensing because the light that is essential for obtaining measurements experiences wavelength-dependent exponential attenuation with distance travelled through the fibre, according to Beer’s Law (also known as the Beer-Lambert law).

Optical fibres function as waveguides due to total internal reflection. Snell’s Law of Refraction describes how light is bent or refracted when it travels through two isotropic, homogeneous media with different indices of refraction $n_1$ and $n_2$:

$$\sin \theta_2 = \frac{n_1}{n_2} \sin \theta_1$$

**EQUATION 1**

where $\theta_1$ and $\theta_2$ are the angles of incidence and refraction, respectively. When light travels from a medium of greater refractive index into a medium with a lower index ($n_1 > n_2$) with an incident angle greater than a critical value $\theta_{crit}$, **EQUATION 1** is no longer valid as it would require $\sin \theta_2$ to be greater than 1 ($\sin \theta_2 > 1$). In this case, the light is not refracted and is instead reflected into the medium having the greater refractive index; this phenomenon is known as total internal reflection. The critical angle $\theta_{crit}$ is given by

$$\theta_{crit} = \sin^{-1} \left( \frac{n_2}{n_1} \sin (\theta_2 = 90°) \right)$$

**EQUATION 2**

$$\theta_{crit} = \sin^{-1} \left( \frac{n_2}{n_1} \right)$$

**EQUATION 3**

When the angle of incidence $\theta_1$ is greater than the critical angle $\theta_{crit}$, light is reflected within the optical fibre core rather than refracted in the cladding.

**FIGURE 3**
Optical fibres are constructed of silica (SiO₂) glass and have two primary components: the core and the cladding. The core and the cladding have different indices of refraction that are chosen to maximize the total internal reflection by minimizing the normal critical angle, allowing light to travel for kilometers with minimal attenuation (FIGURE 3).

The core is therefore the central component of an optical fibre in which most of the light travels. The silica in the core is often doped with germanium dioxide (GeO₂) or aluminium dioxide (Al₂O₃) to slightly increase the refractive index of the core relative to the cladding.

The optical fibres used for distributed sensing are of two main types: multimode or singlemode (FIGURE 4). To mechanically strengthen and protect the fragile silica core/cladding, a primary coating is applied. Use of different materials as the primary coating allow for optical fibre installations from standard to extreme temperature ranges, as listed in FIGURE 4.

<table>
<thead>
<tr>
<th>Primary coating ambient temperature range</th>
</tr>
</thead>
<tbody>
<tr>
<td>Acrylate</td>
</tr>
<tr>
<td>HT Acrylate</td>
</tr>
<tr>
<td>Silicone PFA</td>
</tr>
<tr>
<td>Polymide</td>
</tr>
<tr>
<td>Metal</td>
</tr>
</tbody>
</table>
Typical multimode (MM) fibres have a 50 μm or 62.5 μm diameter core, which allows for significantly more light to travel in the core than in singlemode (SM) fibres and allows for light propagation in a variety of modes.

The refractive index of the core is on the order of 1% greater in magnitude than that of the standard 125 μm diameter cladding (FIGURE 5). The cross-section of most MM fibres also have a graded index (GI), meaning the transition in refractive index is gradual between the core and the cladding.

When an initial pulse of laser light is emitted into an optical fibre, the main mode of light propagation is along the central fibre axis. However, some light may enter the fibre at an angle to its center-line, resulting in internal reflections that cause the light to have a zig-zag or spiral path through the fibre. In this case, the optical path of some of the light rays will be longer than others and will arrive at a specific point along the fibre after the main mode light. GI MM fibres minimize this effect by allowing non-main modes of light to travel faster (due to a slightly lower index of refraction near the edge of the fibre) and better match the arrival time of main mode light at a specific point along a fibre.

Multimode 50/125 μm (core diameter/cladding diameter) GI fibres are most commonly used in DTS, as well as for most high-bandwidth optical fibre communication links. Due to the substantially greater cross-sectional area, the amount of light that can be coupled into and then internally reflected in the core of a MM fibre is higher than in a SM fibre, resulting in a better signal-to-noise (SNR) ratio and resolution performance for MM DTS systems. The greater cross-sectional area minimizes core alignment mismatch resulting in less light lost at splices and mechanical connectors in the optical fibre.

Singlemode optical fibres are characterised by a small core of approximately 9 μm in diameter, which avoids modal dispersion by allowing light to propagate in only one mode. The cladding, like multimode fiber, is most often 125 μm in diameter. The much smaller diameter of singlemode fibres allows much less light to be coupled into the fibre, increasing the challenge of obtaining meaningful measurements from the Raman scattering signal.

Singlemode fibres are most commonly used for distributed acoustic or strain sensing, as these systems rely on the Rayleigh scattering signal, which is orders of magnitude more intense than the Raman scattering signal. Singlemode fibers are typically only used for DTS measurements in retrofit applications where no multimode fiber is available.
A pulsed laser provides the input light source generating the measured backscattered light signal used in DTS.

Though Raman scattering can occur for any wavelength of source light, infrared light is used in fibre optic sensing because the attenuation of light in the fibre is much less at longer wavelengths. The typical wavelengths of light used in DTS are 1064 nm (most applications) or 1550 nm (long ranges) due to favorable attenuation and scattering characteristics at these wavelengths.

Light is emitted by the laser source in short pulses rather than continuously. This allows backscattered light from every point along the fibre to be uniquely identified before another pulse is launched and the scattering location along the fibre to be determined from the time of flight. The length of light pulses and the frequency at which consecutive measurements are taken are crucial in determining the resolution of temperature measurements obtained by a DTS system.

Light travels along the glass core of an optical fibre at a speed \( v \), which is less than \( c \), the speed of light in vacuum (\( c = 3 \times 10^8 \text{ m/s} \)).

This is due to the interaction of the light with the glass, whose refractive index \( n \) is around 1.5 (compared to 1 in vacuum), and is related to \( v \) according to:

\[
\frac{c}{n} \approx 2 \cdot 10^8 \text{ [m/s]}
\]

\( v \) = \frac{c}{n} \approx 2 \cdot 10^8 \text{ [m/s]}

**EQUATION 4**
The duration $\Delta t$ of a light pulse is usually several nanoseconds.

During this time, the light in the fibre, propagating at constant speed $v$ [m/s], covers a distance $\Delta z$ (FIGURE 6) according to:

$$\Delta z [m] = v \left[ \frac{m}{s} \right] \times \Delta t [s]$$

EQUATION 5

For example, a pulse of $\Delta t=10$ ns, travelling at $v=2\times10^8$ m/s, has a distance extension of $\Delta z=2$ m.

In a DTS system, backscattered light returns to the source point, where it is detected, along the same path as the incident light and with a very similar velocity to that of the incident light pulse. An incident pulse of time duration $\Delta t$ (i.e. distance extent $\Delta z$) generates a backscatter signal that is effectively the average backscatter light of a distance extent $\Delta z/2$.

The factor 1/2 comes from the fact that the incident pulse and backscatter light travel in opposite directions. The resulting temperature information is an average over the distance extent of the backscatter light signal. Therefore, the incident pulse width has a direct impact on the spatial resolution of the DTS.

<table>
<thead>
<tr>
<th>Incident pulse time duration $\Delta t$ [ns]</th>
<th>Incident pulse distance extent $\Delta z$ [m]</th>
<th>Backscatter distance extent $\Delta z/2$ [m]</th>
</tr>
</thead>
<tbody>
<tr>
<td>1.25</td>
<td>0.25</td>
<td>0.125</td>
</tr>
<tr>
<td>2.5</td>
<td>0.5</td>
<td>0.25</td>
</tr>
<tr>
<td>5</td>
<td>1</td>
<td>0.5</td>
</tr>
<tr>
<td>10</td>
<td>2</td>
<td>1</td>
</tr>
</tbody>
</table>

**TABLE 1**

Typical pulse durations $\Delta t$ for DTS systems, with corresponding lengths $\Delta z$ and sampling resolutions $\Delta z/2$. Velocity $v$ of light in glass is assumed to be $\sim2\times10^8$ [m/s].
The location corresponding to a measured backscattered light signal is determined from the principles of optical time domain reflectometry.

The distance \( z \) from the DTS interrogator to the back scattered light signal source is determined by measuring the time \( t \) between the launch of a laser pulse at time \( t_0 \) and the detection of the backscattered light at time \( t_1 \):

\[
\frac{z}{2} = v \times t
\]

**EQUATION 6**

Where the speed \( v \) of the traveling light pulse is given by **EQUATION 3** \((-2 \times 10^8 \text{ [m/s]})\) and \( t = t_1 - t_0 \).

Therefore, a light signal detected with a measured travel \( t=100 \text{ ns} \) is correlated with a source 10 m away from the DTS interrogator, while a light signal detected at time \( t=10 \text{ µs} \) correlates to a source 1 km away.

The determination of the distance to an optical signal source by timing the interval between emission of an optical signal and the detection of the backscattered or reflected signal is the basic principle of optical time domain reflectometry (OTDR).
Laser light input into an optical fibre is scattered through three main processes: Rayleigh, Raman and Brillouin scattering.

When light travels through an optical fibre, most of the photons are scattered elastically by silica molecules in the fibre that are much smaller than the wavelength of light in a process called Rayleigh scattering.

In Rayleigh scattering no energy is exchanged between the photons and the scattering molecules, so the wavelength of the scattered light is the same as that of the incident light. A small proportion of incident photons, however, are inelastically scattered through the processes of Raman and Brillouin scattering.

Raman scattering is due to interactions between incident photons and vibrational molecular states that result in excitation events. Unlike Rayleigh scattering, Raman scattered photons have wave properties that differ significantly from the incident light. If an electron excited by an incident photon relaxes into an energy state greater than its initial state, a photon with less energy than the incident light is emitted (the molecule gains energy from the interaction).

If the molecule loses energy in the collision however, the scattered photon will have a higher energy (shorter wavelength) than the incident photon. These two possible outcomes of an interaction result in two spectral features in the backscattered light, referred to as the Stokes ($S$) and anti-Stokes ($aS$) bands, with wavelengths above and below that of the incident light (FIGURE 7).
The Stokes and anti-Stokes bands are symmetrical about the wavelength of the Rayleigh scattered light because they correspond to the energy difference between the same two upper and lower energy states within the scattering molecule (FIGURE 8). In optical fibre glass the typical wavelength shift of the Raman bands is ±~50 nm around the Rayleigh peak, which coincides with the wavelength of the emitting laser source (1064 or 1550 nm in DTS systems).

The intensities of the Stokes and anti-Stokes signals differ because the intensity depends on the relative populations of energy levels in the fibre molecules, which is temperature dependent. At thermal equilibrium, most of the molecules in a material are at a lower energy state and will absorb energy from incident photons, leading to more transitions that result in Stokes emissions than anti-Stokes emissions. The intensity of the anti-Stokes signal is therefore highly temperature dependent and the Stokes scattering signal is typically more intense.

In addition to the Rayleigh and Raman backscatter signals, a third backscattered signal, Brillouin scattering, is due to the interaction between photons and physical variations in the fibre material. Brillouin backscatter signal is the main effect that DSTS systems rely on, as this scattering mechanism is mediated by the index of refraction of the material, which, for clear materials, changes when the material experiences deformation. Material deformation varies with temperature, so Brillouin scattering is somewhat temperature dependent.
Incident photons may gain or lose energy when interacting with variations in the material, so as in Raman scattering, both wavelength upshifted and downshifted Brillouin bands are observed. Rayleigh scattering can also depend on variations in the scattering material, but such scattering is due to random fluctuations rather than the correlated fluctuations that cause Brillouin scattering. The wavelengths of the two Brillouin bands vary similarly with temperature and strain.

The wavelengths of laser light used in DTS applications are selected due to the attenuation of light in optical fibre. Light travelling in optical fibre experiences exponential attenuation with distance according to the Beer-Lambert Law, which relates the gradual loss of light with distance traveled to the refractive properties of the transmitting material. Attenuation is wavelength dependent, and the Stokes and anti-Stokes signals therefore experience different losses (FIGURE 9).

The wavelengths of emitted laser light in DTS systems are therefore chosen for their relatively favorable positions in the spectrum of attenuation of light in an optical fibre. 1064 nm is in the middle of a monotonic and almost linearly decreasing section of the attenuation spectrum, whereas 1550 nm corresponds to the absolute minimum of attenuation. The greater attenuation around 1064 nm (particularly for the anti-Stokes signal) makes DTS systems emitting light at this wavelength preferable for shorter length fibres, while systems with laser sources emitting at 1550 nm are mainly used for long-range applications.
The temperature at sample locations along an optical fiber is determined from the measured Raman backscattered light signals and key calibration parameters.

The intensity of the Raman backscattered signal, composed of Stokes and anti-Stokes signals, due to light emitted through an optical fiber in a DTS system depends on the populations of quantum states within the fiber molecules.

The population of quantum states is temperature dependent, so temperature can be derived from the measured Raman scattering signal. The intensity of the photons backscattered in the anti-Stokes band increases with higher temperature, and vice versa. The Stokes signal is much less sensitive to temperature and is used for reference to quantify the variations in anti-Stokes backscatter along the optical fiber. The temperature $T \text{[K]}$ at distance $z \text{[m]}$ along the physical fiber is therefore determined from the ratio of the intensities of the Stokes $I_S$ and anti-Stokes $I_{aS}$ signals.

A widely accepted (Dakin et al., 1985) formulation of the mathematical relation is:

$$ T(z) = \frac{\gamma}{K - \Delta \alpha \cdot z - \ln\left(\frac{I_S(z)}{I_{aS}(z)}\right)} $$

(EQUATION 7)

where

- $\gamma \text{[K]}$ depends on the distribution of quantum states and is a system constant fixed by manufacturing. $\gamma$ quantifies the shift in energy between the photons travelling at the wavelength of the incident light and the backscattered Raman photons. ($\gamma = h \Omega / k$ with $h$ being Planck’s constant, $\Omega$ is the difference in frequency between the incoming laser pulse and the backscattered Stokes radiation, and $k$ is Boltzmann’s constant.)

- $K$ is a dimensionless calibration parameter related to the temperature offset that accounts for the efficiency of the DTS components which are common factors for all the measured traces (e.g. laser optical power, efficiency of the photodiodes converting the backscattered light into electrical signal, electronic circuits).

- $\Delta \alpha \text{[m}^{-1}\text{]}$ is the differential loss calibration parameter that accounts for the difference in attenuation between the Stokes and anti-Stokes signal due to their differing wavelengths. $\Delta \alpha$ is often expressed in [dB/km].

- $z \text{[km]}$ is the fibre distance. $z = 0$ at the DTS instrument, though a distance offset can be specified when configuring data collection.

- $I_S(z)$ and $I_{aS}(z)$ are the Stokes and anti-Stokes (respectively) signal amplitudes over distance.

$\gamma$, $K$ and $\Delta \alpha$ must be determined through DTS system calibration in order to obtain accurate distributed temperatures from EQUATION 7.
Temperature, temporal and spatial resolutions are the key performance factors of a DTS system.

DTS systems provide the spatial distribution of temperature along a length of optical fibre in time, so temperature and spatial resolution, as well as the time required to achieve high resolution measurements, are all key factors to consider in DTS applications.

These factors are inextricably linked and depend on the system pulse width, pulse repetition rate, sampling resolution and the length of fibre. The interplay of performance factors is an important consideration in the design of any DTS system and is critical for insuring quality and reliability in DTS measurements as well as for determining the optimal DTS system and settings for a given application.

The sampling resolution of a DTS system describes its ability to convert a true continuous spatial distribution of temperature into discrete measurements.

The sampling resolution of a DTS system is the smallest length increment a DTS system can sense (or sample) over the entire length of an optical fibre; The DTS system provides one averaged temperature measurement per sample.

The sampling resolution of a DTS system is determined by the sampling frequency of the DAQ (data acquisition) card which is typically implemented with a FPGA (field-programmable gate array) and specialized high-speed A/D (analog to digital converters) chip technology.

Although the spatial resolution and sampling resolution are related, they must not be confused with each other. The sampling resolution cannot be equal to the spatial resolution.

To satisfy the Nyquist sampling criterion, the sampling rate must be more than twice the highest frequency component of a signal to properly capture the signal. Similarly, the spatial resolution cannot be smaller than two consecutive samples.

In general, the spatial resolution is slightly larger than two times the sampling resolution. Oversampling at much greater than one half the spatial resolution results in increased data volumes without significant additional information contained within the data-set.
A DTS system’s ability to resolve temperature changes along an optical fibre is described by the spatial resolution.

Each temperature measurement provided by a DTS system is averaged over a specified length increment, so the sensor output response to a change in temperature along the fibre is somewhat blurred at the edges of the change. The spatial resolution of a DTS system is therefore determined by applying a step change in temperature between two adjacent lengths of fibre (10 m or more) and determining the distance needed to capture between 10% and 90% of the variation (FIGURE 10). Typically, a temperature step of about 30°C is applied. The 10%-90% definition of spatial resolution is appropriate for determining the degree to which a transition can be reproduced in the sensor output.

**FIGURE 10**
DTS temperature with distance data from a spatial resolution test for the Silixa Ultima-S (sampling resolution = 0.125m). The distance required to capture between 10% and 90% of the applied temperature step change is about 30 cm, indicating the spatial resolution of this system.
It is important to note that different DTS manufacturers may apply different definitions, e.g. defining the sampling resolution as the spatial resolution, varying the amplitude of the step change in temperature, and/or accounting for another percentage of detection (e.g. between 20% and 80%). Extra care should be taken in reading the specifications of a DTS system to be aware of the definition of spatial resolution that is being applied.

The temperature distribution indicated by DTS data is the convolution of the sensor’s spatial response with the true spatial distribution of temperature. If the temperature features are finer than the spatial resolution of the sensor, the details of temperature variations along the fibre will be unresolvable.

If there is a hot spot (or a cold spot) along the fibre with an extent smaller than the spatial resolution, the temperature corresponding to the center of the hot spot will be lower than the actual value (FIGURE 11). The spatial resolution, together with the range, determines the number of independently resolvable points that can be measured by the system.

**FIGURE 11**
In this experimental set-up to determine the spatial resolution of a DTS, multiple fibre optic sections of different lengths (left to right: 10m, 4m, 2m, 1m, 0.5m, and 0.25m) are exposed to a fixed temperature step change. The only hot spot which cannot be fully resolved above the 90% step change corresponds to the 0.25m length of fiber which is shorter than the 30cm spatial resolution of the Ultima-S DTS unit.
The measurement range determines the pulse repetition rate of the laser.

The optical interrogator in a DTS system must not trigger a light pulse until the previously emitted pulse has travelled to the end of the optical fibre and the back-scattered light has returned to the detector from every scattering point along the entire length of the fiber. If pulses are emitted any faster, the backscatter from each pulse will overlap at the detector and the data cannot be correlated with location along the fibre. A DTS system’s range is therefore the maximum length of fibre that can be measured, which dictates the maximum pulse repetition rate.

This is mathematically shown by assuming that the emitted laser light and the backscattered Stokes and anti-Stokes bands all travel at the same speed $v$ such that the minimum time $\Delta t$ between triggers can be approximately determined as

$$\Delta t = 2 \frac{\Delta z}{v} = 2 \frac{L \cdot n}{c}$$

(EQUATION 8)

where $L$ is the total length of the fibre and the speed of light in the fiber $v = c/n$ (EQUATION 3). This leads to a launch repetition rate $R$:

$$R = \frac{1}{\Delta t} = \frac{c}{2 \cdot L \cdot n} \propto \frac{1}{L}$$

(EQUATION 9)

which is inversely proportional to the maximum optical fibre length to measure. For example, if $L=5$ km and $n=1.5$, the minimum time between triggers would be $\Delta t=5 \mu s$, and therefore $R=20$ kHz (TABLE 2).

<table>
<thead>
<tr>
<th>Range (km)</th>
<th>Max pulse rep rate (kHz)</th>
</tr>
</thead>
<tbody>
<tr>
<td>5</td>
<td>20</td>
</tr>
<tr>
<td>10</td>
<td>10</td>
</tr>
<tr>
<td>15</td>
<td>6.7</td>
</tr>
<tr>
<td>20</td>
<td>5</td>
</tr>
<tr>
<td>25</td>
<td>4</td>
</tr>
<tr>
<td>30</td>
<td>3.3</td>
</tr>
<tr>
<td>35</td>
<td>2.9</td>
</tr>
</tbody>
</table>

(TABLE 2: Maximum pulse repetition rate in kHz for different maximum measurement ranges.)

The temperature resolution of a DTS system is the smallest temperature variation a system can detect or resolve.

The temperature resolution (e.g. 0.01 °C or 0.1 °C) and is an indicator of the precision, or repeatability, of a temperature measurement along the fibre. Temperature resolution degrades exponentially along a fibre because the backscattered signal used to determine temperature is increasingly attenuated with fibre distance.

The temperature resolution $\sigma_T(x)$ of a DTS system along a length $x$ of optical fibre is assessed as the standard deviation of the $N$ consecutive samples defining the section $x$:

$$\sigma_T(x) = \sqrt{\frac{1}{N-1} \sum_{i=1}^{N} (T(i)_x - \overline{T(x)})^2}$$

(EQUATION 10)
where \( T(x) \) is the temperature array and \( \overline{T(x)} = \frac{1}{N} \sum_{i=1}^{N} T(i) \) is the average temperature for each considered section \( x \).

The temperature resolution for a section \( x \) is typically defined by considering \( N=10 \) consecutive spatial samples. The temperature resolutions of consecutive sections \( x \) are then calculated by applying a moving window along the fibre. So for a measurement with 25 cm spatial sampling, the rolling window, centered on the measurement location, is 2.5 m long; in the case of 1 m spatial sampling, the rolling window is 10 m long.

For \textit{Equation 10} to reflect the true temperature resolution, it is assumed that the temperature is constant within the rolling window so that only the noise variation is taken into account. This is an important consideration, particularly when analysing field data where real temperature features along the fibre could mask the temperature resolution estimation by adding unwanted feature variation to the noise variation estimation.

If a more precise estimate of resolution is needed, a fitting function is used to determine the final temperature resolution along the fibre. The temperature resolution degrades exponentially over distance as the optical signal propagates following the Beer-Lambert Law. Therefore, a fitting function of the form

\[
\sigma_T(z) = a \cdot e^{bz}
\]

\textit{Equation 11}

is used to estimate the final resolution function with distance (\( z \) is the fibre position, \( a \) and \( b \) are fitting parameters). In \textit{Figure 12}, an example of \( \sigma_{T(x)} \) is shown in red and the exponential fit \( \sigma_{T(z)} \) in blue.

The launch repetition rate, or the number of pulses a DTS system can launch per second, has a significant impact on the temperature resolution, as well as the other performance factors. As the repetition rate increases, so does the number of samples that can be averaged per second. The measurement time defines the time required to obtain readings for all points along the sensing fibre to achieve the required temperature resolution.
The temperature resolution of a DTS system is primarily limited by the exponential attenuation of the signal with distance traveled in the fibre and the Signal-to-Noise Ratio (SNR) of the system. In general, the backscatter signal in a DTS system from an individual light pulse is quite small, so the SNR is very low. Given a fixed signal level, the SNR and temperature resolution can be improved by averaging multiple measurements. Averaging can be done both in time (longer averaging time) and space (increased sampling resolution).

The optoelectronic noise for each given data point can be characterized over time as additive white Gaussian noise (AWGN).

Therefore, averaging generates a reduction of the noise floor equal to the square root of the number of samples averaged. This has a direct translation into the temperature resolution of the DTS. The temperature resolution of a DTS system improves by averaging more measurements in time (FIGURE 13).

Generally, from a known temperature resolution $\sigma_m$ of a measurement of duration $m$, the temperature resolution $\sigma_n$ of a measurement of any duration $n$ (not necessarily a multiple of $m$) can be found as:

$$\sigma_n = \frac{\sigma_m}{\sqrt{n/m}}$$

(EQUATION 12)

FIGURE 13

Temperature resolution curves for a DTS system sampling every 0.25m, at different measurement times. The optical fibre (length = 5000m) was coiled and kept at constant controlled conditions in a laboratory temperature chamber.
Space-domain averaging can provide similar temperature resolution improvements as time-domain averaging.

The trade-off in this case is that the spatial resolution is broadened. The spatial averaging improvement factor is similar to the time-domain averaging improvement as long as the original sampling interval is above the spatial resolution of the DTS. This is because consecutive temperature samples are only statistically uncorrelated if the original sampling interval is larger than the spatial resolution. For example, for a DTS system with 30cm spatial resolution, the improvement factor when averaging consecutive 50cm samples will be \( \frac{1}{\sqrt{2}} \), but will be much smaller when averaging consecutive 12.5cm samples.

A change in the maximum fibre length (range) results in a change in the maximum pulse repetition rate, which affects the number of raw measurements per second.

The subsequent change in temperature resolution follows similarly from Equation 11. The temperature resolution at a fixed spatial sample along the fibre, for different maximum fibre lengths \( L_1 \) and \( L_2 \) is therefore given by:

\[
\sigma_{L_1} = \frac{\sigma_{L_2}}{\sqrt{\frac{L_2}{L_1}}}
\]

Equation 13

As the fibre length increases, total attenuation increases with distance and subsequently the temperature resolution degrades. The temperature resolution is degraded because as fibre length increases, the maximum possible pulse repetition frequency decreases.

For a fixed measurement time, fewer samples are collected and averaged. Use of longer width pulses and wider sampling intervals however can result in acceptable temperature resolution for measurements using longer fibres.
3.5 Accuracy

Accuracy refers to the agreement between the temperature measured by a DTS system and the true temperature of the optical fibre.

The terminology often used for distributed systems varies from that of conventional point transducers. In properly configured DTS systems, accuracy is constant with fibre distance and is defined as the difference of the average value of the measured temperature profile from the actual temperature, while temperature resolution is the spread in the average temperatures measured (FIGURE 14). Accuracy depends upon numerous parameters affecting the measurements, including the environmental conditions of the acquisition unit and the stability of the measurements. Typically, the measurement resolution defines the smallest temperature step the unit can measure and is mainly affected by the optical pulse energy, noise in the photoreceivers, and analogue electronic front-end bandwidth. The magnitude of both the accuracy and resolution contribute to the measurement uncertainty or error of the measurements.
There are a variety of DTS system configuration options available which affect how temperature offset and differential attenuation calibration parameters are accounted for when determining temperature with distance.

Accurate DTS temperature measurements at a distance $z$ (m) depend on the ability of the instrument to measure the intensities of the backscattered Stokes and anti-Stokes signals, which is impacted by physical parameters of the DTS system: $\gamma$, $K$, and $\Delta\alpha$. The determination of these three parameters defines the precision and accuracy of a DTS system's temperature measurements and can be explicitly calculated or resolved, depending on the configuration of the DTS system. Implicit determination of these calibration parameters is not the norm, but, for some applications, doing so can provide additional configuration flexibility and allow for improved temperature resolution.

DTS systems can function with a variety of configurations of optical fibre, which affect how the system resolves temperature. DTS units are designed to allow for data collection from single or multiple fibres, which are connected to various emission/detection channels. The two most basic configurations are referred to as single-ended and double-ended (FIGURE 15). In a single-ended configuration, laser pulses are input and backscattered light is measured at the same end of the fibre, which is connected to one channel of the DTS. In a double-ended configuration, the fibre extends from the DTS and is looped back so that each end of the fibre is connected to differing detection channels. Laser pulses are input and monitored from both ends of the fibre (or both channels), alternating sides during each measurement interval.

Explicit calculation of $\gamma$, $K$, and $\Delta\alpha$ for any system configuration relies on three independent comparisons of three observed reference temperatures, usually with two as close as possible to the laser source where light attenuation with distance is negligible (to calibrate $\gamma$ and $K$, respectively) and one at the far end of the optical fibre (to calibrate $\Delta\alpha$). However, in practice, $\gamma$ is a fixed value determined by the instrument manufacturer and has proven to be insensitive to the main factors affecting the other parameters, so it is reasonable to assume that it is constant in time. Therefore, for most applications, the differential attenuation $\Delta\alpha$ and temperature offset $K$ parameters are the parameters that must be determined, and they are derived differently depending on the DTS system configuration. Calibration procedures can be greatly simplified by the capabilities of a given DTS unit and the use of specific cable configurations which can reduce the number of external, independent temperature comparisons needed and the amount of data post-processing required, without a decrease in accuracy.
In addition to the losses experienced by light in a fibre due to wavelength dependent attenuation, bends, connectors, splices, and other irregularities or defects can all contribute to losses in transmitted and detected light signals in DTS systems.

Attenuation caused by these various fibre features may also vary with fibre temperature, resulting in temporal variations in attenuation. In the case that attenuation is uniform along a fibre, \( \Delta \alpha \) can be determined through calibration based on independent temperature measurements at two points along the cable and the temperature can be determined from Equation 6, assuming all other parameters are known. \( \Delta \alpha \) can also be set at a fixed, constant value (e.g. 0.1 dB/km) according to the specifications of the optical fibre being used. In most DTS applications however, attenuation varies along the cable due to combinations of some of the irregularities mentioned. In this case, all attenuating effects between the DTS unit at \( z' = 0 \) and the sample location \( z' = z \) must be combined. In this case, the temperature equation takes the form:

\[
T(z) = \gamma \frac{Y}{K - \int_0^z \Delta \alpha(z')dz'} - \ln \left( \frac{k(\gamma)}{\alpha_0(\gamma)} \right)
\]

EQUATION 14

The integral \( \int_0^z \Delta \alpha(z')dz' \) has the unit of dB, so must be scaled by \( \ln(10)/10000 \) to change the units from dB to linear.

Calibration measurements necessary to determine differential attenuation in DTS systems can be achieved by holding a section or, ideally, two sections of optical fibre at known, constant temperatures. Typically, sections of optical fibre are coiled and placed into temperature-controlled baths inside insulated containers such that the fibre does not contact the container and experience temperature variations.

Since temperature is averaged over a length of fibre, the section of fibre in the bath should be at least 10-25 m long and bath temperature should be continuously and independently monitored for reference; a 100 Ohm platinum resistance thermistor (Pt100) is the standard sensor used. If a fibre is composed of several joined sections of fibre, each section must be individually calibrated.

Most DTS systems include automated calibration algorithms for correcting for the differential attenuation using different system configurations. These so-called differential loss corrections terms are known as fixed value, temperature matching, external probe, or double-ended. Fixed value, temperature matching, and external probe options are most commonly used for single-ended fibre configurations.

In a DTS system with a single-ended configuration, both the temperature offset and differential attenuation parameters must be determined through calibration unless a differential loss correction term can be specified in the DTS system configuration. Proper differential loss correction is required to maintain accuracy with distance from the DTS unit for single-ended measurements.

\( T(z) = \gamma \frac{Y}{K - \int_0^z \Delta \alpha(z')dz'} - \ln \left( \frac{k(\gamma)}{\alpha_0(\gamma)} \right) \)

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In a DTS system with a single-ended configuration, both the temperature offset and differential attenuation parameters must be determined through calibration unless a differential loss correction term can be specified in the DTS system configuration. Proper differential loss correction is required to maintain accuracy with distance from the DTS unit for single-ended measurements.
The fixed value differential loss correction term is the difference in attenuation experienced by the Stokes and anti-Stokes signals expressed in [dB/km].

Usually the optimum setting can be found by ensuring any points along the trace, which are known to have the same temperature, line up correctly. If no such points exist, then the default values of 0.250 [dB/km] for Silixa DTS systems with 1064 nm wavelength lasers or 0.020 [dB/km] for systems with 1550 nm wavelength lasers should be used. The DTS system applies the chosen fixed value differential loss correction term to all data sets.

Temperature matching may be chosen for determining the differential loss correction term when the temperature along two separate sections of the fibre is, on average, the same.

This option is usually only chosen when there is symmetry in the fibre path; for example, when the fibre is deployed to the far-end of a linear structure and brought back over part of the return path in a so-called “U-shaped” or “J-shaped” configuration.

When temperature matching is chosen, the positions and extent of two sections that have the same average temperature must be defined. For this technique to work optimally, each section should be as long as possible, i.e. the user-defined fibre sections should contain as many data points as possible, and as far apart as possible. Typically, each section should be at least 500 m long, with the nearest section in the first quarter of the fibre length and the furthest section in the last quarter of the fibre. If temperature matching is used, the differential loss is re-calibrated each time a measurement is taken.

The external probe differential loss correction can be used when it is possible to place two external temperature probes at a position with the same temperature as the average temperature along a section of fibre.

This option is only accurate if the section is in the most distant two-thirds of the fibre length, so, in practice, the external probe option is generally only used when the end of the fibre returns to a position very close to the DTS unit. When using the external probe option, the position and extent of the reference section must be defined and the differential loss is re-calibrated each time a measurement is made.
In the common case of non-uniform attenuation, it is appropriate to use a double-ended configuration which allows for a uniform level of precision along the entire length of the fibre.

To determine the double-ended differential loss correction term, laser pulses are launched into the fibre from alternating ends of the fibre (e.g., a pulse is launched in the ‘forward’ direction from the end of fibre connected to Channel 1 during the first measurement period and launched in the ‘reverse’ direction from Channel 2 during the second measurement period, and so on). Following Dakin et al. (1985), the ‘forward’ laser pulse starts at $z=0$ and the ‘reverse’ pulse starts at $z=L$, with $L$ being the length of the cable. If the temperature over a span $\Delta z$ remains constant during the forward and reverse measurement periods, EQUATION 14 can be used to solve for the differential attenuation $\Delta \alpha$ by setting $T(z)=T(L-z)$ and rearranging:

$$\int_{z}^{z+\Delta z} \Delta \alpha (z') dz' = \ln \left( \frac{l_5 (z+\Delta z)}{l_{aS} (z+\Delta z)} \right) - \ln \left( \frac{l_5 (z)}{l_{aS} (z)} \right) + \ln \left( \frac{l_5 (z+\Delta z)}{l_{aS} (z+\Delta z)} \right) - \ln \left( \frac{l_5 (z)}{l_{aS} (z)} \right)$$

EQUATION 15

The arrows ($\leftarrow, \rightarrow$) indicate whether the measurement is considered in the forward or reverse direction. By obtaining a measurement from each sample $\Delta z$ along the entire length of the fibre and integrating the values, the differential attenuation $\Delta \alpha$ at each location $z$ can be determined. The data used is a time-average over the entire measurement period, during which it is assumed $\Delta \alpha$ remains constant.

Longer measurement integration times and spatial averaging are used so that the value of differential attenuation obtained through calibration does not add significantly to the uncertainty of DTS temperature measurements. Since the temperature reported for each sample in a double-ended system is determined from signals travelling from both directions in the fibre, the measurements are noisiest closest to the sensor and have the least noise at the midpoint of the fibre.

Double-ended calibration uniquely and robustly corrects for differential light loss that varies over the length of the fibre. There is no requirement for the fibre to follow the same path or for fibre at a given distance from the unit on both sides of the loop to have the same temperature, though double-ended calculations do require very precise alignment of the forward and reverse signals.

Using a loop of fiber with double-ended differential loss calibration on a linear object requires the length of fibre to be doubled. Also, the interrogator must probe twice the length of an equivalent single-ended system, which may lead to decreased resolution for longer lengths of fibre.

Use of this method is mostly recommended when the fibre is located in a particularly harsh environment and is likely subject to damage-induced changes in differential loss. In other cases, it is generally advisable not to use double-ended calibration because of the complexity of the configuration required.
FIGURE 17
Optical power exponential attenuation of Stokes and anti-Stokes signals for 10 km of optical fibre placed in a temperature chamber at 30°C. The attenuation of the emitting source would lie in between the two plotted signals. The emission wavelength is 1064 nm; at 1550 nm, the slope differential between the attenuation curves of Stokes and anti-Stokes is much smaller, because of differing attenuation coefficients.
4.1.2 Determination of temperature offset

The temperature offset accounts for all optoelectronic conversion and efficiency factors that affect the Raman backscattered signals (filters, amplifiers, attenuators, etc.).

The temperature offset therefore has the same effect on each temperature sample (FIGURE 18), unlike the differential attenuation, which generally has a greater effect on samples farther from the DTS unit (FIGURE 16). Some of the quantities characterizing the temperature offset are sensitive to the operational temperature changes experienced by a DTS instrument which may cause the offset to vary in time.

Unless the DTS instrument can be maintained at a very stable temperature, the temperature offset must be determined for each measurement. An uncalibrated temperature offset manifests as a distant-invariant deviation from the true temperature for each sample. While this does not affect the accuracy of the measured relative changes in temperature, it does affect the accuracy of the absolute temperature measurements.

From the known value of $\gamma$ and the determined value of the differential attenuation $\Delta\alpha$, the value of the temperature offset is easily obtained using a known temperature reference to calculate an offset for the entire data set. Typically, the temperature offset is dynamically determined within the DTS instrument using a highly accurate internal temperature reference; however, for unique applications, offset calibration using a probe and/or an external fiber reference section may be appropriate.

Use of a reference section and a probe is preferable over a single point-source measurement because this method allows for the noise in the measurement to be averaged out and improves the temperature resolution.

An external probe to determine the temperature offset can be used when the probe can be placed such that it is at the same temperature as the average temperature along a selected section of the sensing fibre. In practice, the external probe option is only used to determine the temperature offset when an insulated coil of the sensing fibre is positioned, along with the probe, near the DTS unit. When this option is chosen, the DTS ignores the internal reference and instead uses the selected external probe and reference section as the inputs into the temperature offset calculation. If the external probe option is used for both the differential loss...
4.1.2 Determination of temperature offset

and temperature offset calibrations, then the reference section and probe must be at different locations or temperatures in both cases.

Typically, the reference section consists of several meters of coiled optical fibre placed in a homogenous hot or cold water bath, or in a thermally isolated container. To achieve accurate calibration, the fibre temperature must be constant in space and preferably constant, or slowly varying in time, during the DTS system measurement to assure that the temperature measured by the external probe accuracy reflects that of the reference fibre.

The use of water mixers is highly recommended during calibration to avoid the generation of vertical temperature gradients within the bath. To minimally reduce the effects of heat exchange between the environment and the bath, the coil of reference cable should be placed horizontally within the bath and contact should be avoided with the walls.

The length of the reference section of cable used for calibration is not fixed, though the maximum possible optical fibre length that can be placed in the bath while following the stated guidelines is recommended.

Longer reference fibres provide more data points, and thus a more statistically robust average temperature measurement for comparison to the independent probe measurement. For example, if 10 m of optical fibre is used as a reference section with a sampling resolution of 12.5 cm, the average temperature is determined from 80 samples.

The difference between the temperatures measured simultaneously by the independent probe and by the reference fibre determines the offset, which is then added (with the appropriate sign) to all points of the DTS measurement. The offset should be calculated and corrected for each DTS measurement.

4.1.3 Explicit mathematical determination of physical parameters

The explicit determination of $\gamma$, $K$, and $\Delta\alpha$ via three independent temperatures is the most complicated to realize, as it requires access to the optical fibre both in proximity and far from the DTS instrument, accurate external temperature measurements, and significant post-processing of data.

While internal configuration options are often suitable for calibration, explicit determination of parameters may be necessary for DTS systems for which an estimation of internal unit parameters is unavailable or a simple single-ended fibre optic configuration is used. For some applications doing so can also provide additional configuration flexibility allowing for improved temperature determination.

A linear system of three temperature equations must be solved iteratively for each measurement to determine temperature. If three independent temperature measurements are unavailable for comparison, optimization methods must be used to determine the values of the unknown parameters ($\gamma$, $K$, and $\Delta\alpha$).
In active distributed sensing, a thermal signal is input into the environment along an optical fibre path and monitored in time by a DTS system.

Active DTS methods are of particular use in hydrogeological and atmospheric DTS applications in which environmental properties can be determined from the measured temperature profile in time. Here we will briefly discuss only the method which makes use of a hybrid (composite) cable containing both an optical fibre used for DTS measurement and a heating element with a constant power output per meter.

The amount of thermal transport away from the cable can be determined from the measured variations in temperature along the fibre. The measured temperature response in time caused by environmental heat transport can be used to determine the thermal properties of the measurement environment, including the apparent thermal conductivity of material surrounding the fibre. Active DTS also provides a way to characterize a variety of physical properties that affect thermal conductivity, such as seepage, soil moisture, or even wind speed.

Active heating is much more common than active cooling because an electrical resistance circuit can easily be used as a linear heater in the cable. Cooling methods such as circulation of fluids or utilising refrigeration are far more challenging.

![FIGURE 19](image)

Example active DTS system setup using a DTS unit and heat pulse control unit with a composite cable containing both optical fibre and conductive wire.
The heating system in active DTS usually consists of a simple circuit with a hybrid fibre/electrical conductor cable used as a resistor with a constant electrical resistance per meter voltage is applied across the conductor(s) in the cable to induce current flow.

If the resistance and voltage are held constant, the power output per meter of cable is unchanging in time (FIGURE 19). A constant power output along the cable is generally required for calculating parameters from the temperature response due to the inherent variability of many environmental factors, such as fluid flow (FIGURE 20).

Pulse modulation of heating can be used in some cases to characterize the contrast between heated and unheated environmental conditions. The design of an active DTS system is dependent on the physical quantities that are to be derived from the measured temperature response.
Cable manufacturers provide the electrical resistance per unit length (Ω/m) of the conducting materials embedded in hybrid measurement cables.

Ohm’s law,

\[ I = \frac{V}{R} \]

**EQUATION 16**

where \( I \) is the current in amperes (A) through the conductor, \( V \) is the voltage applied across the conductor in volts (V), and \( R \) is the resistance in ohms (Ω), can therefore be used to determine the voltage necessary to achieve the desired power rate per unit length in Watts per meter (W/m), where

\[ \text{Power} = V \times I = I^2 \times R = \frac{V^2}{R} \]

**EQUATION 17**

Typical values can vary from 1-2 W/m for atmospheric measurements such as wind speed, to 10-20 W/m for geothermal wells and shallow soil monitoring (FIGURE 21).

The thermal properties of the environment surrounding the heated cable and the resolution of the DTS system dictate the amount of voltage necessary to apply to the conductor to achieve the desired power output and to generate a clearly detectable thermal signature.

![FIGURE 21](Image)

Power output of a cable with 2x18 AWG (American Wire Gauge) electrical conductors each with electrical resistance of 21.4 Ohms/km for different voltages as a function of cable length. Voltage is applied to one end of the cable with the two conductors connected in series.
### 5.3 Hardware for active DTS

The components required for active DTS measurements typically include:

**DTS**
Due to the rapid temperature changes caused by active heating, a DTS unit capable of high temporal resolution is desired. Additionally, fine temperature resolution will allow for small temperature changes to be readily quantified, which can improve the sensitivity of the active DTS system and allow for the use of lower heat outputs per unit length of cable.

**Hybrid (composite) fibre optic cable**
Hybrid or composite cables integrate electrical conductors into the fibre optic cable to allow a single cable to be used for heating and measuring the temperature response. Cables should be optimized for active DTS measurements when possible and material choice carefully considered with temperatures during heating often approaching the melting point of common cable materials.

**Heating control system**
A system that integrates with the DTS and provides constant power output in time is desired.

**Power source**
Generally, a ~240 V, 30 A power source is required, though the exact requirements are dependent on the cable design, measurement length, and desired power output per unit length.

### 5.4 Advantages & challenges of active DTS

The main advantage of using an active DTS system is the ability to measure properties and processes related to heat transport when there is no reliable natural signal that can be used.

By creating a known and controlled thermal stress and monitoring the environmentally regulated response, active DTS can be used to characterize many properties and processes that would not be feasible otherwise.

Active DTS poses challenges that are not associated with passive DTS methods. Additional equipment such as a heating control system and specialized cables are required. The power consumption of the active DTS system may be excessive for long term remote monitoring where there is no main power and generators can't be left running.

Also, the maximum measurement range can be limited by the power available at the installation. However, with proper system design challenges associated with active DTS measurements can usually be overcome.
Light travelling through an optical fibre experiences an exponential attenuation of the intensity $I$ with distance $z$.

Attenuation is a physical process independent of the resolved temperature and leads to an unavoidable degradation of temperature resolution with distance. Beer's law describes this fundamental intrinsic fibre loss as an exponential decay with distance $z$ of the initial (i.e. $z=0$) light intensity $I_0$ according to:

\[
I(z) = I_0 e^{-\alpha z}
\]

**EQUATION 18**

where $\alpha$ [1/m] is the wavelength-dependent attenuation coefficient of the travelling wave and $I_0$ is the intensity at $z=0$. Alternatively, $\alpha$ is also often expressed in [dB/km], which is related to the [1/m] value by the following relation:

\[
\alpha \text{ [dB/km]} = \left( \frac{10^4}{\ln(10)} \right) \cdot \alpha \text{ [1/m]}
\]

**EQUATION 19**

For the temperature determination at a certain distance $z$ along the optical fibre it is necessary to calculate the intensity $I_S(z)$ and $I_{aS}(z)$ of the Stokes and anti-Stokes backscatters, respectively. $I_S(z)$ and $I_{aS}(z)$ are expressed as:

\[
I_S(z) = C_S e^{-\alpha R z} e^{-\alpha S z} \left\langle n_k \right\rangle + 1
\]

**EQUATION 20**

\[
I_{aS}(z) = C_{aS} e^{-\alpha R z} e^{-\alpha S z} \left\langle n_k \right\rangle
\]

**EQUATION 21**

where

- $C_S$ and $C_{aS}$ are constants that account for all the components of a DTS system that are common factors for the measured traces; in particular, the efficiency of the receivers and of the amplifier system.
- $e^{-\alpha R z}$ = primary pulse attenuation, common for anti-Stokes and Stokes.
- $e^{-\alpha S z}$ = Stokes attenuation
- $e^{-\alpha aS z}$ = anti-Stokes attenuation

\[
\left\langle n_k \right\rangle + 1 = \frac{e^{-\hbar \Omega / k T(z)}}{1 - e^{-\hbar \Omega / k T(z)}} + 1 = \text{Stokes quantum mechanical term}
\]

\[
\left\langle n_k \right\rangle = \frac{e^{-\hbar \Omega / k T(z)}}{1 - e^{-\hbar \Omega / k T(z)}} = \text{anti-Stokes quantum mechanical term}
\]

where

- $\hbar = h/2\pi$
- $\Omega$ = angular frequency [rad/s] = Red (>1) or Blue (<1 ) frequency shift
- $h = 6.626 \cdot 10^{-34}$ [m² kg/s] = Planck's constant
- $k = 1.38 \cdot 10^{-23}$ [J/K]
- $E=\hbar \nu=\hbar \Omega$ = Photon energy [J] [J/K]
- $\frac{\hbar \Omega}{k}$ = non-system dependent constant ≈ 600 [K]
\( I_{AS}(z) \) and \( I_{S}(z) \) account for the fact that the primary pulse intensity (at \( z=0 \)), is attenuated at a rate \( \alpha_R \) of the incident light (same as the Rayleigh wave) at \( z \), whereas the returning signals to \( z=0 \) experience an attenuation \( \alpha_{AS} \) for the anti-Stokes band and \( \alpha_{S} \) for the Stokes band.

The attenuation terms \( \alpha_{AS}, \alpha_R \) and \( \alpha_{S} \) differ because the wavelengths of the emitting laser source and backscattered anti-Stokes and Stokes signals are different and therefore experience different losses.

The Stokes and anti-Stokes quantum mechanical terms both contain a temperature dependence, though the anti-Stokes signal proves to have greater temperature dependency. Therefore, the Stokes band is therefore not completely insensitive to temperature variations. However, it is used as a reference and the DTS temperature information is obtained from the ratio of the Stokes and anti-Stokes intensities, as this allows simplifying the calculations:

\[
\frac{I_{S}(z)}{I_{AS}(z)} = \frac{C_{S}}{C_{AS}} e^{-\Delta \alpha \cdot z} e^{-\frac{h\Omega}{kT(z)}}
\]

\text{EQUATION 22}

where \( \Delta \alpha = \alpha_{S}-\alpha_{AS} = \) differential attenuation [m\(^{-1}\)] or \( \left[ \frac{dB}{km} \right] > 0 \).

Solving the above equation for the temperature (in Kelvin) results in:

\[
T(z) = \frac{\gamma}{K - \Delta \alpha \cdot z - \ln \left( \frac{l_{S}(z)}{l_{AS}(z)} \right)}
\]

\text{EQUATION 23}

where \( \gamma = \frac{h\Omega}{K} \) and \( K = \ln \left( \frac{C_{S}}{C_{AS}} \right) \).
### Glossary

6.1

This list of DTS terminology is intended to be a reference guide for basic definitions of DTS terms.

<table>
<thead>
<tr>
<th>Term</th>
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<tbody>
<tr>
<td><strong>Accuracy</strong></td>
<td>Quantifies how close a measurement is to a real value. For DTS systems, the real values are absolute temperatures. Accuracy must not be confused with resolution or precision.</td>
</tr>
<tr>
<td><strong>Active DTS</strong></td>
<td>A method of DTS data collection based on using a heated fibre optic cable(s) and monitoring the thermal response during the heating and/or the cooling phase. Also commonly known as the heat pulse method. Active DTS can be used for measurement of thermal properties and directly dependent quantities (e.g. soil moisture, heat fluxes), and of advective fluxes of water and air.</td>
</tr>
<tr>
<td><strong>Anti-Stokes Backscatter</strong></td>
<td>The component of Raman backscattering that is shifted to a lower wavelength from the incident light. This backscattered signal is strongly temperature dependent.</td>
</tr>
<tr>
<td><strong>Attenuation</strong></td>
<td>A measure of the loss of intensity with distance of the light travelling in an optical fibre, usually expressed in dB or dB/km.</td>
</tr>
<tr>
<td><strong>Backscattered Light</strong></td>
<td>Also referred to as backscatter, is the portion of light that, during the forward propagation of a pulse through the optical fibre, is scattered back towards the light-emitting source because of collisions of the photons with the constitutive elements of the optical fibre.</td>
</tr>
<tr>
<td><strong>Brillouin Scattering (Brillouin Effect)</strong></td>
<td>Light that is shifted in wavelength from an incident light pulse through energization of the silica lattice structure. The Brillouin backscatter caused by an incident light pulse propagating through an optical fibre is most commonly used for distributed strain sensing. Brillouin wavelengths are shifted less from the incident wavelength than the Raman backscatter.</td>
</tr>
<tr>
<td><strong>Calibration</strong></td>
<td>When used in conjunction with DTS this term refers to means of calibrating the physical quantities that allow inferring temperature from DTS raw signal. A good calibration is crucial to avoid the DTS temperature readings being affected mainly by offset (constant shift between measured and real absolute temperatures) and slope issues (variation with distance of DTS readings of a constant temperature). Multiple calibration methods are available including both single-ended and double-ended techniques.</td>
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### Glossary

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<tr>
<td>Cladding</td>
<td>Outer structure of an optical fibre in direct contact with the fibre core with a lower refractive index than the core causing light to be contained within the fibre through total internal reflection.</td>
</tr>
<tr>
<td>Core</td>
<td>The central component of an optical fibre in which the majority of light travels. This material has a higher refractive index than the cladding in order to maintain total internal reflection.</td>
</tr>
<tr>
<td>Differential Attenuation</td>
<td>The difference between the loss of light intensity with distance for the Stokes and anti-Stokes Raman backscatter. This physical quantity requires careful calibration in order to set the correct slope of a DTS trace (DTS readings of a constant temperature do not change with distance).</td>
</tr>
<tr>
<td>Distributed Acoustic Sensor (DAS)</td>
<td>A sensor that allows for the collection of acoustic signal fully distributed along an optical fibre. The most advanced DAS systems allow for the collection of quantitative coherent full-waveform acoustic data.</td>
</tr>
<tr>
<td>Distributed Temperature Sensor (DTS)</td>
<td>A sensor that allows for the collection of temperature data fully distributed along an optical fibre.</td>
</tr>
<tr>
<td>Double-ended</td>
<td>A method of DTS measurement and calibration that requires both ends of the forward optical path to be connected to the DTS instrument. Each temperature trace is provided by combining the forward and reverse raw data traces measured sequentially from each end. Particularly useful for DTS calibration of fibre optic cables in harsh environments or with multiple splices and/or tight bends as spatially variable differential attenuation is compensated for.</td>
</tr>
<tr>
<td>Duplexed</td>
<td>A configuration where two fibres in the same fibre optic cable are spliced at the far end to realize a U-shaped configuration. Duplexing allows performing measurements using either one DTS optical channel with one end of the optical fibre connected, or two DTS optical channels with each end of the optical fiber connected. In the latter case, measurements are taken sequentially from the two ends and can be single-ended (forward/reverse raw data traces are not combined), or double-ended.</td>
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<td><strong>Fibre in Metal Tube (FIMT)</strong></td>
<td>A typically stainless-steel tube, often gel-filled that incorporates optical fibres in a hermetically sealed assembly. Commonly used for high temperature and harsh environment cabling. Sometimes referred to as tube encapsulated fibre (TEF).</td>
</tr>
<tr>
<td><strong>Fibre Optic Cable</strong></td>
<td>A complete assembled product including optical fibre(s), strength members, and protective jacket(s) that protect the optical fibre(s) from environmental factors.</td>
</tr>
<tr>
<td><strong>Fusion Splicing</strong></td>
<td>A method of joining two optical fibres end to end using intense heat. This method can provide connections with very low attenuation.</td>
</tr>
<tr>
<td><strong>Heatable Fibre Optic Cable</strong></td>
<td>Fibre optic cable including metal elements (usually copper wires) close to the fibres optics core. Electrical current sent through them allows generating a distributed heat pulse along the cable. Heatable cables are used for the Active DTS method. Sometimes are also called hybrid cables or composite cables.</td>
</tr>
<tr>
<td><strong>Incident Light</strong></td>
<td>The light that is coupled into an optical fibre at the DTS and travels down the fibre from the near end toward the far end of the fibre.</td>
</tr>
<tr>
<td><strong>Multi-mode Fibre</strong></td>
<td>Optical fibre that has a relatively large core diameter (when compared with single-mode fibre) of usually 50 or 62.5 µm. This type of fibre supports more than one light propagation mode.</td>
</tr>
<tr>
<td><strong>Optical Fibre</strong></td>
<td>Consisting of the core surrounded by the cladding, the optical fibre provides a means of transmitting signal at very high bandwidths over long distances through a process known as total internal reflection. Optical fibres are usually made of silica.</td>
</tr>
<tr>
<td><strong>Optical path</strong></td>
<td>Indicates the path travelled by the light in the optical fibre(s). When forward optical path is intended the path is in the direction of the emitted light. The end of the forward optical path can be either the far end of the fibre optic cable or the near if a duplex configuration is applied.</td>
</tr>
<tr>
<td><strong>Passive DTS</strong></td>
<td>The standard method of collecting temperature data along an optical fibre using DTS instrumentation, without artificially heating the fibre optic cable as in the active DTS methods.</td>
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<td><strong>Precision</strong></td>
<td>Quantifies the smallest variation of a quantity that an instrument is able to detect and properly measure above the noise. For a DTS system, the smallest fluctuations in temperature that can be resolved. Precision can be used as synonymous of resolution, and is an index of the repeatability (in both space and time) of a measurement. It must not be confused with the accuracy. An instrument can be very precise, but still poorly accurate, and vice-versa.</td>
</tr>
<tr>
<td><strong>Raman Scattering (Raman Effect)</strong></td>
<td>Light that is shifted in wavelength from an incident light pulse through vibrations of the silica molecular structure. The Raman backscatter caused by an incident light pulse propagating through an optical fibre is most commonly used for DTS. Raman wavelengths are shifted further from the incident wavelength that the Brillouin backscatter.</td>
</tr>
<tr>
<td><strong>Rayleigh Scattering (Rayleigh Effect)</strong></td>
<td>Light scattered at the same wavelength of the incident light. In an optical fibre Rayleigh scattering is caused by microscopic density and refractive index variations. The Rayleigh backscatter caused by an incident light pulse propagating through an optical fibre is most commonly used for DAS.</td>
</tr>
<tr>
<td><strong>Repeatability</strong></td>
<td>Refers to the precision (or resolution) at which a measurement can be repeated in space and time. It must not be related to the accuracy. For a DTS system, the index of repeatability is provided by the temporal resolution.</td>
</tr>
<tr>
<td><strong>Sampling Resolution</strong></td>
<td>The sampling interval spacing along an optical fibre that a DAS or DTS system collects. Though related, the sampling resolution should not be confused with the spatial resolution of the instrument.</td>
</tr>
<tr>
<td><strong>Single-ended</strong></td>
<td>A method of measurement and calibration in which temperature traces are collected from one end of an optical fibre.</td>
</tr>
<tr>
<td><strong>Single-mode Fibre</strong></td>
<td>Optical fibre that has a relatively small core diameter (when compared with multi-mode fibre) of usually 8 or 9 µm. This type of fibre supports only a single light propagation mode.</td>
</tr>
<tr>
<td><strong>Spatial Resolution</strong></td>
<td>The minimum distance that a distributed sensing system needs in order to measure a step change in signal along an optical fibre. For DTS systems, this is usually calculated as the distance needed to differentiate between 10% and 90% of an applied temperature step.</td>
</tr>
<tr>
<td><strong>Step loss</strong></td>
<td>Sharp loss in intensity of light propagation along an optical fibre, very localized in space. It is usually due to connectors and fusion splices, and appears as a step in the Stokes and Anti-Stokes traces. If the step loss experienced by the Stokes and anti-Stokes optical signals is different, this can cause a step artefact on the temperature trace which appears from the location of the step loss until the end of the sensing fibre. Double-ended configurations typically automatically account for step losses.</td>
</tr>
<tr>
<td><strong>Stokes Backscatter</strong></td>
<td>The component of Raman backscattering that is shifted to a higher wavelength from the incident light. This backscatter is not strongly temperature dependent.</td>
</tr>
<tr>
<td><strong>Temperature Resolution</strong></td>
<td>Indicates the precision of the DTS system to measure temperature (i.e. how noisy is the instrument). It is usually calculated in space as the standard deviation of DTS measurements along short optical fibre sections (in order to not consider differential attenuation effects) held at constant temperature. It can also be calculated in time for a fixed point along the optical fibre.</td>
</tr>
<tr>
<td><strong>Temporal Resolution</strong></td>
<td>Indicates the time required to collect an individual trace along an optical fibre.</td>
</tr>
<tr>
<td><strong>Termination</strong></td>
<td>The completion at either end of an optical fibre or cable. Examples include connectors and U-bends.</td>
</tr>
<tr>
<td><strong>Tube Encapsulated Fibre (TEF)</strong></td>
<td>See Fibre in Metal Tube (FIMT).</td>
</tr>
<tr>
<td><strong>U-bend</strong></td>
<td>Usually indicates the bending of 180 degrees of two fibres spliced at the far end of a fibre optic cable. This doubles the forward optical path within the same cable and allows an easy application of the methods requiring a U-shaped configuration.</td>
</tr>
<tr>
<td><strong>U-shaped configuration</strong></td>
<td>Deployment of a fibre optic cable with both ends of the forward optical path connected to a DTS instrument. Realized by u-bending either two fibres optics (at the far end) or the cable itself (at half of its length). This configuration provides temperature traces mirrored at the U-bend. Moreover, it allows performing measurements in two directions. It is necessary for the double-ended method.</td>
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Smolen, J.J., and A. van der Spek, *Distributed Temperature Sensing: A DTS Primer for Oil & Gas Production* (May 2003), Shell International Exploration.

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